

# Quantum Noise in Gravitational Wave Detectors

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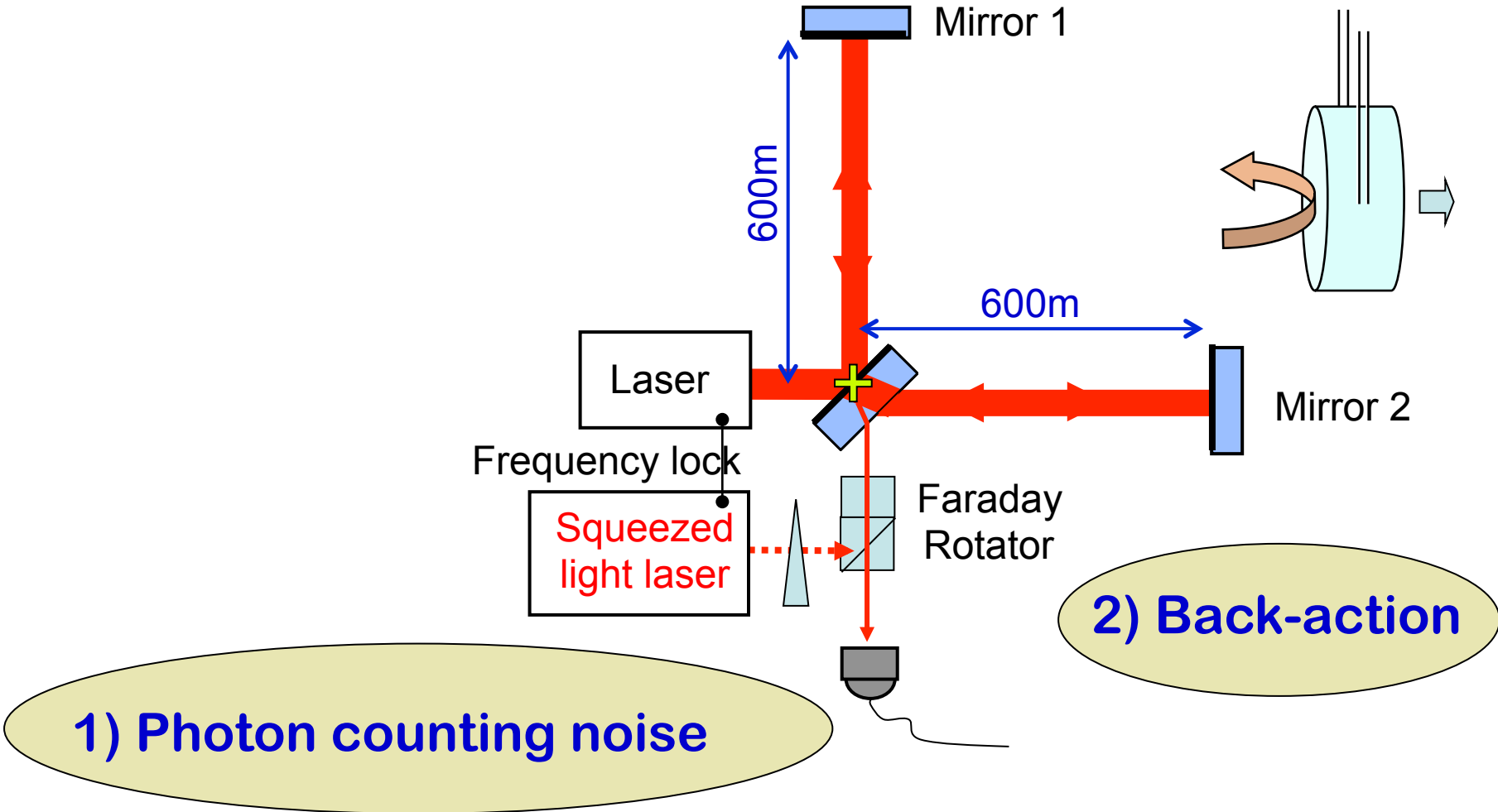
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Institut für Gravitationsphysik

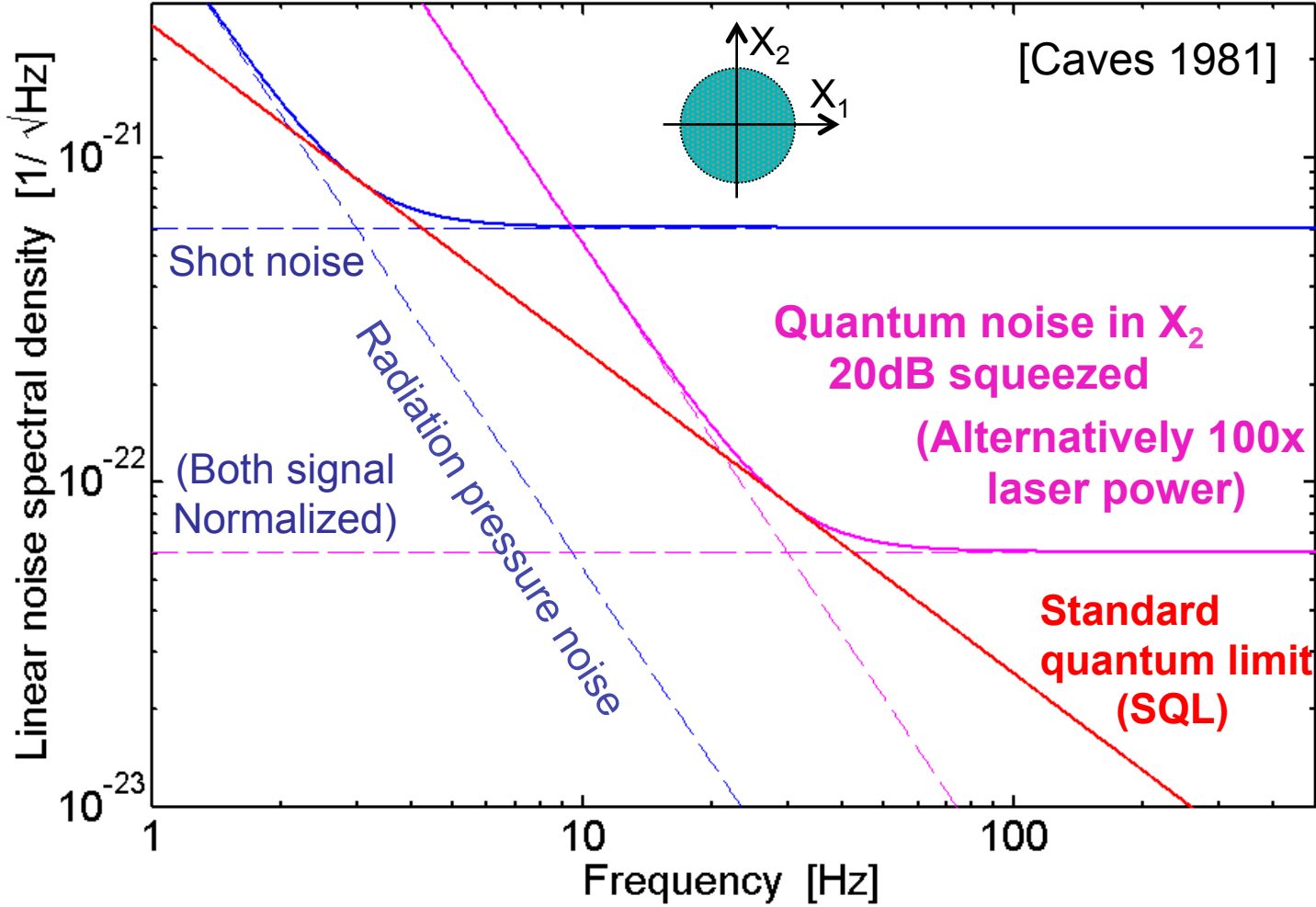
Leibniz Universität Hannover



# Michelson Laser Interferometer



# Measurement Noise and Back-Action



# The SQL – an ultimate Limit?

## PHYSICAL REVIEW LETTERS

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### **Defense of the Standard Quantum Limit for Free-Mass Position**

Carlton M. Caves

*Theoretical Astrophysics, California Institute of Technology, Pasadena, California 91125*

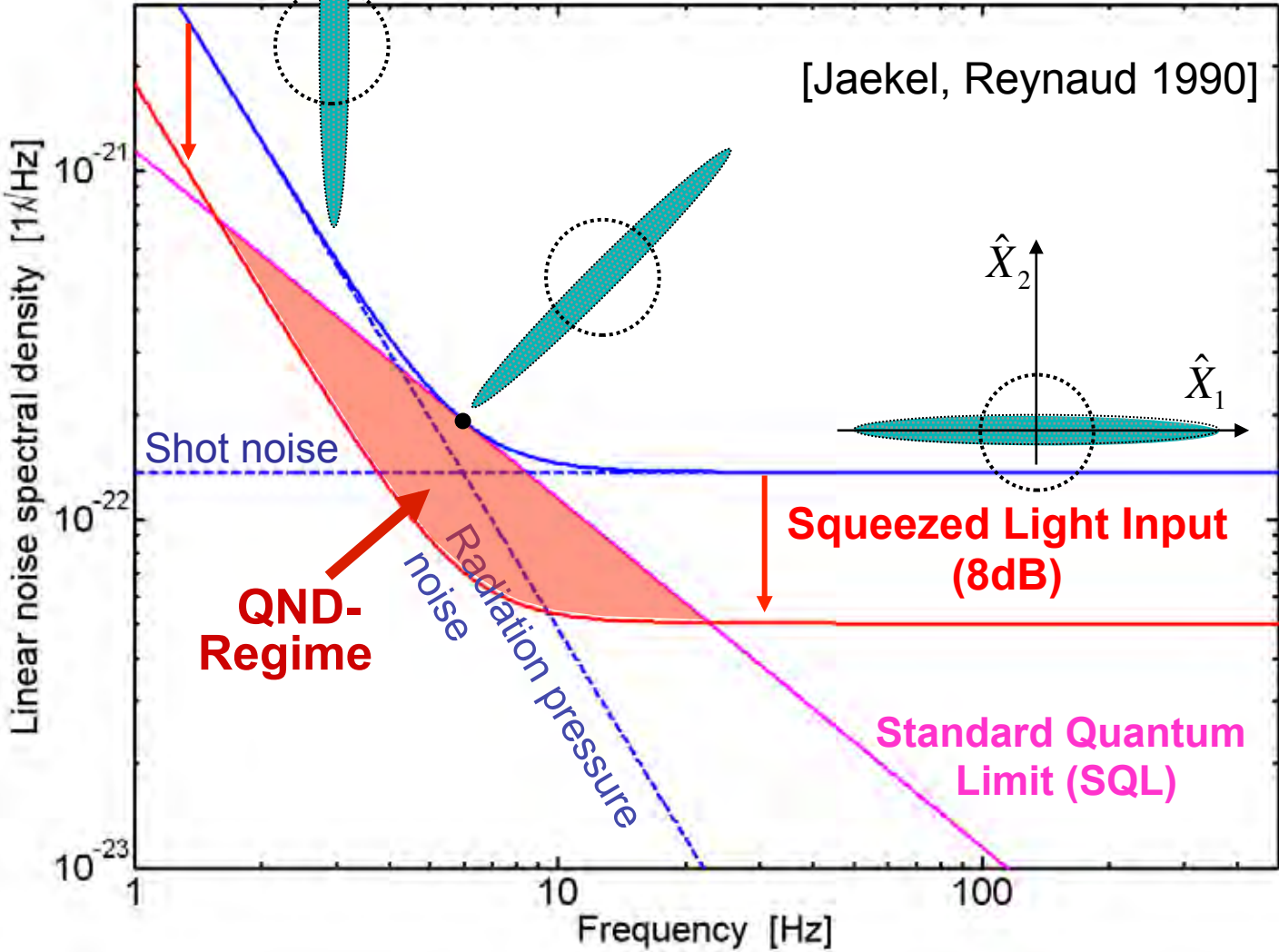
(Received 6 April 1984)

Measurements of the position  $x$  of a free mass  $m$  are thought to be governed by the standard quantum limit (SQL): In two successive measurements of  $x$  spaced a time  $\tau$  apart, the result of the second measurement cannot be predicted with uncertainty smaller than  $(\hbar\tau/m)^{1/2}$ . Yuen has suggested that there might be ways to beat the SQL. Here I give an improved formulation of the SQL, and I argue for, but do not prove, its validity.

[C. M. Caves, *Phys. Rev. Lett.* **54**, 2465 (1985)]



# Squeezing SN and RPN



*Europhys. Lett.*, 13 (4), pp. 301-306 (1990)

## Quantum Limits in Interferometric Measurements.

M. T. JAEKEL(\*) and S. REYNAUD(\*\*)

(\*) *Laboratoire de Physique Théorique de l'Ecole Normale Supérieure*(§)  
*24 rue Lhomond, F-75231 Paris Cedex 05*

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*4 place Jussieu, F-75252 Paris Cedex 05*

“Photon counting noise and radiation pressure noise in a GW detector can **both** be squeezed!”

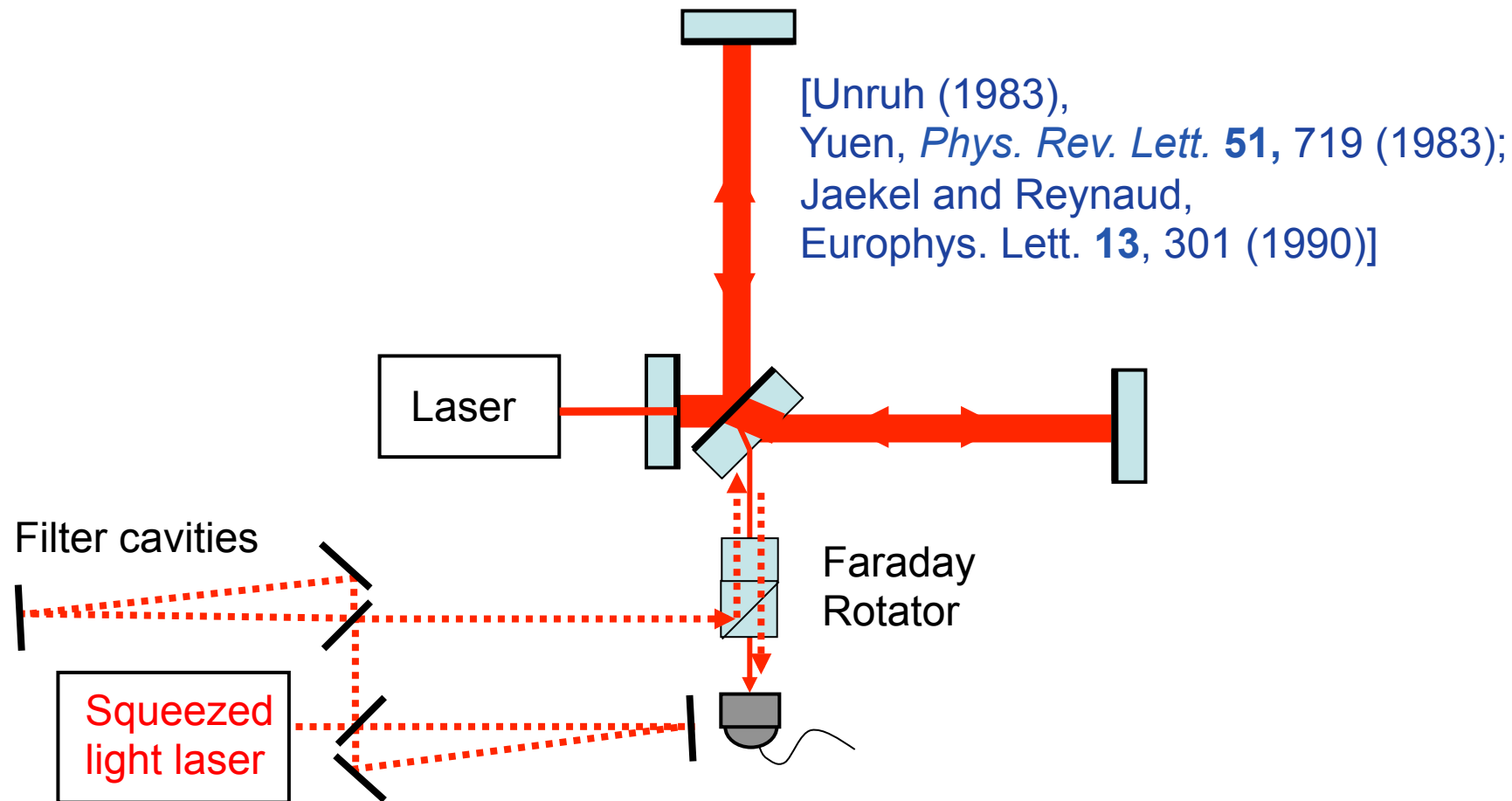
First doubts on the free-mass SQL as the ultimate limit:

W. G. Unruh, in *Quantum Optics, Experimental Gravitation, and Measurement Theory* 647–660 (Plenum, 1983),

H. P. Yuen, *Phys. Rev. Lett.* **51**, 719 (1983)



# Squeezing SN *and* RPN: Filter cavities

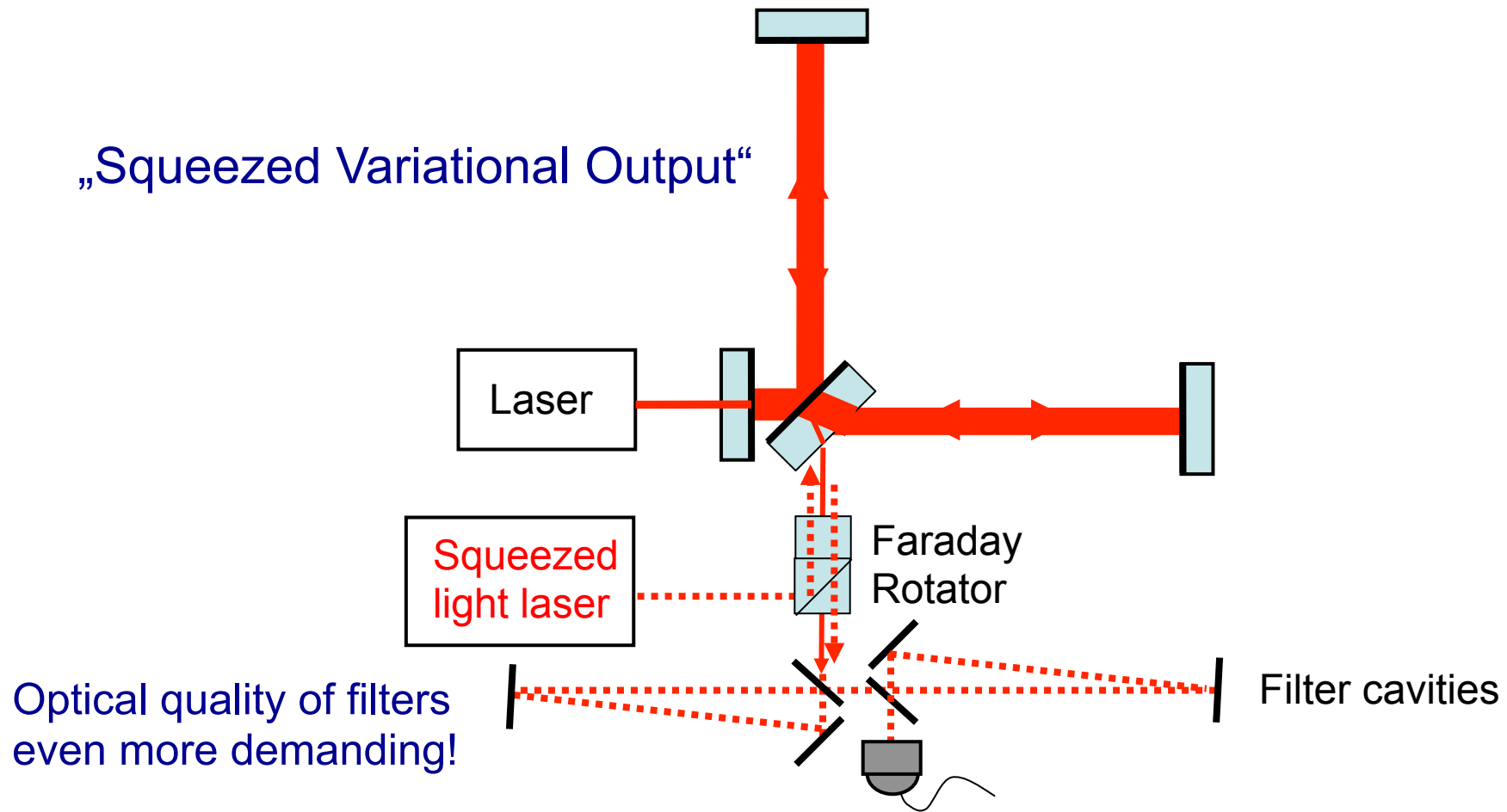


[Unruh (1983),  
Yuen, *Phys. Rev. Lett.* **51**, 719 (1983);  
Jaekel and Reynaud,  
*Europhys. Lett.* **13**, 301 (1990)]

[Kimble *et al.*, *Phys. Rev. D* **65**, 022002 (2001)]



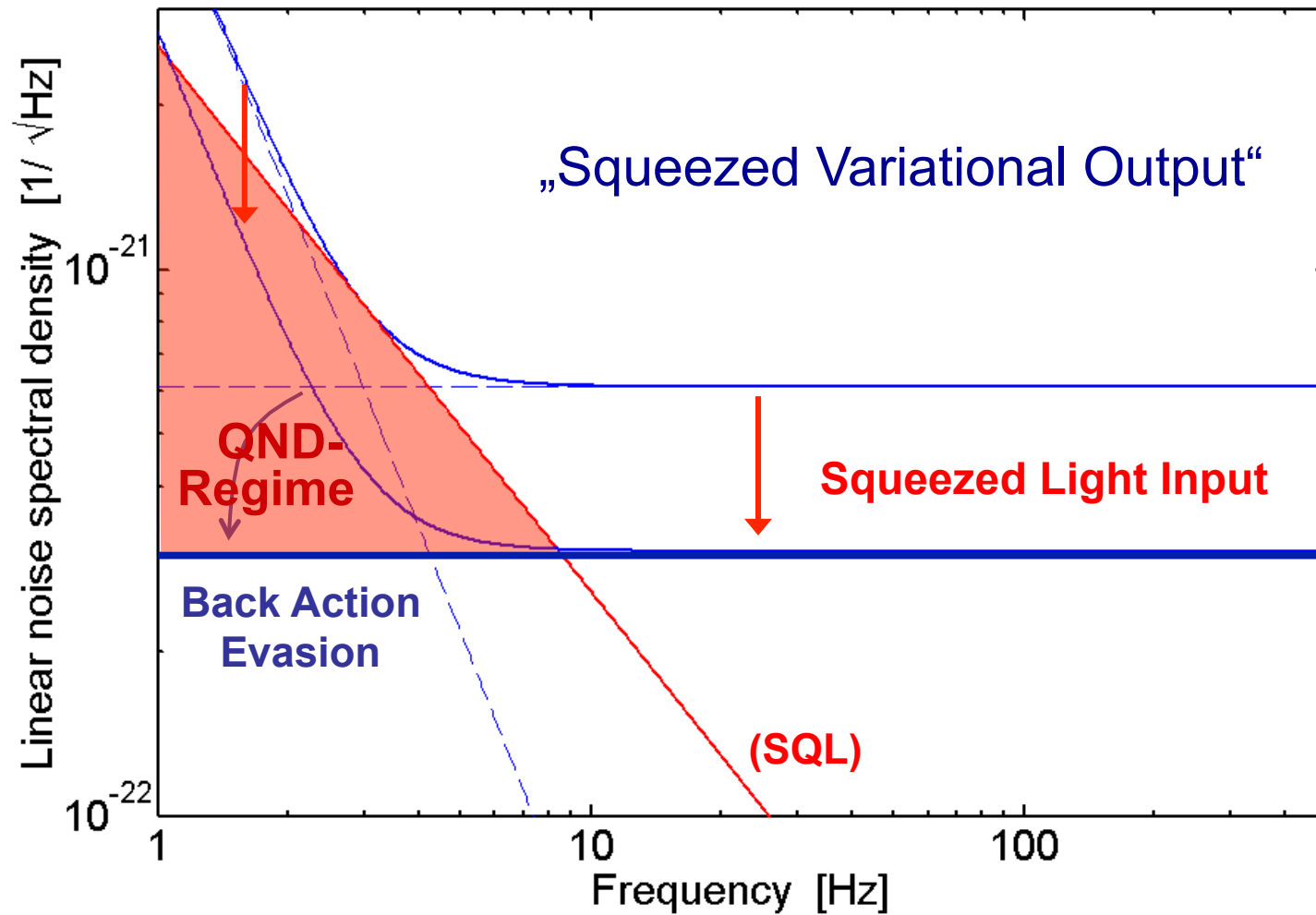
# Squeezing and Full Evasion of RPN



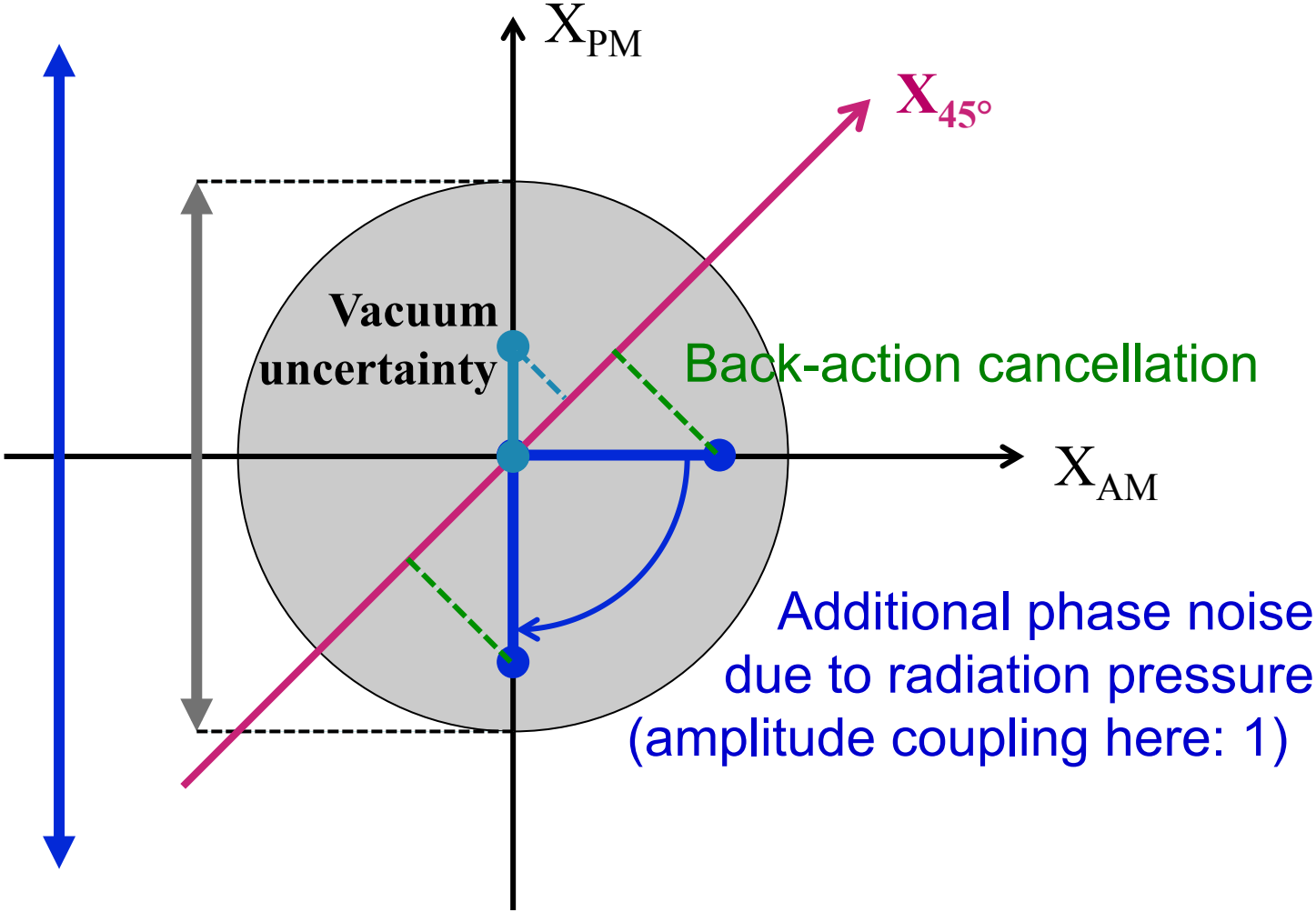
[Kimble *et al.*, Phys. Rev. D **65**, 022002 (2001)]



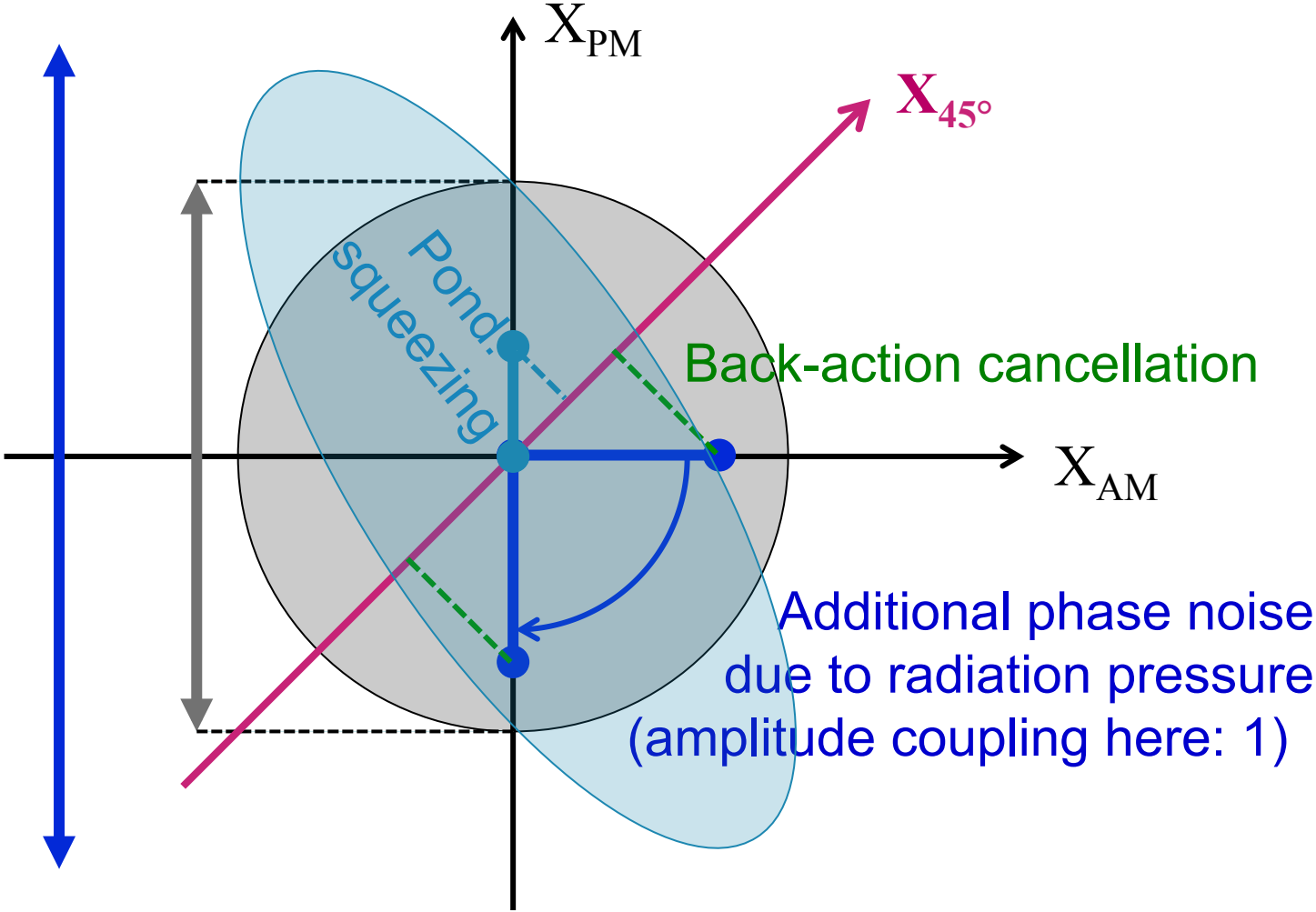
# Squeezing and Full Evasion of RPN



# Ponderomotive Squeezing



# Ponderomotive Squeezing



# Ponderomotive Squeezing

- Is of rather limited strength (requires RPN being dominant)
- has a frequency dependent angle (and strength)
- good to cancel back-action, less good as a squeezing source

## Recent experiments:

- Observation of radiation pressure noise on a membrane at NIST  
[T. P. Purdy, R. W. Peterson, C. A. Regal, Science **339**, 801 (2013)]
- Observation of ponderomotive squeezing at Caltech  
[A.H. Safavi-Naeini, S. Gröblacher, J.T. Hill, J. Chan, M. Aspelmeyer, O. Painter, arXiv:1302.6179]



# Quantum noise in GW detectors

18/06/17

- 1) Photon counting noise (no real mirror displacement)
- 2) Radiation pressure noise  
(momentum transfer from light to mirror)
- (3) zero-point fluctuation of the mirror  
is not an issue! )

"Opto-mechanics does not need  
an optical cavity"



Again: We are interested in Modulations  
of the light power!

- 1) Photon counting noise has a white spectrum!  
" Shot noise <sup>(SN)</sup> per  $\Delta\Omega = 1\text{Hz}$  equivalent to a  
variance of apparent mirror displacement  $\Delta^2 x$   
in  $\text{m}^2$  for a simple Michelson interferometer

$$\Delta^2 x_{SN} = \frac{h c \cdot \lambda}{4\pi P_0} \equiv S_{X,SN} \quad \text{in} \quad \left[ \frac{\text{m}^2}{\text{Hz}} \right]$$

" Displacement normalized (single-sided)  
power spectral density "

## 2) Radiation pressure noise

Force :  $|\vec{F}| = \frac{2 \cdot P_0}{c}$  ← retro-reflection  
← light power reflected

Quantum uncertainty

$$\Delta^2 P_0 = 2 t \omega P_0 \equiv S_{P_0, SN} \quad \text{in} \quad \left[ \frac{W^2}{Hz} \right]$$

$$\Rightarrow \left[ S_{F, SN} = \left( \frac{2}{c} \right)^2 \cdot S_{P_0, SN} = \frac{8 t \omega P_0}{c^2} \quad \text{in} \quad \left[ \frac{N^2}{Hz} \right] \right]$$

Independent of Fourier frequency  $\Omega$   
 $\leq$  white spectrum.

The effect of the back-action force on mirror displacement (rpm)

- effect is stronger the smaller the mass
- effect strong on mechanical resonance

$$S_{x, rpm} = |H(\Omega)|^2 \cdot S_{F, SW}$$

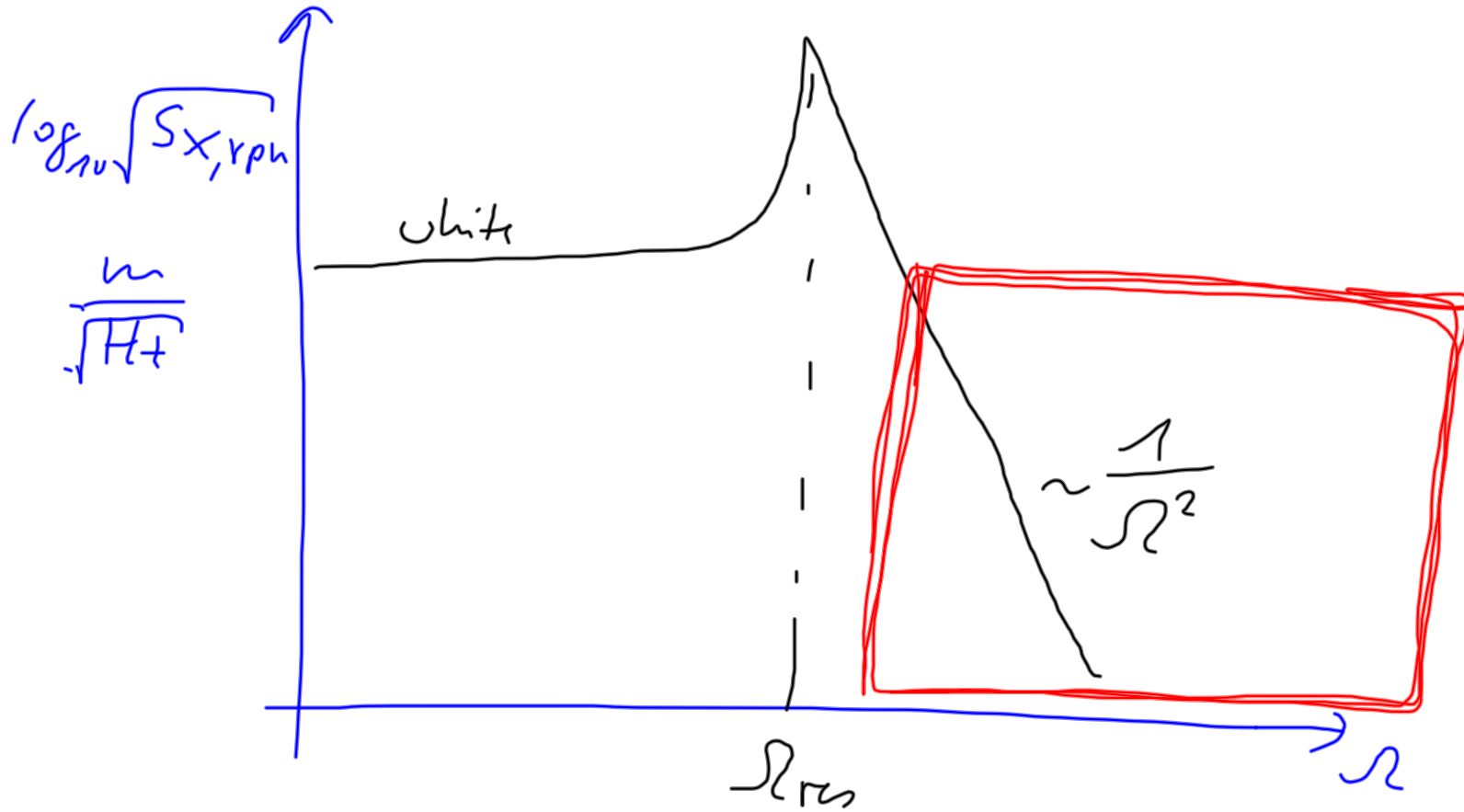
$$= \left| \frac{-1}{m(\Omega^2 + i\gamma\Omega - \Omega_{res}^2)} \right|^2 \cdot \frac{8t\omega P_0}{c}$$

"Radiation pressure noise"

$$\text{in } \frac{m^2}{Hz}$$



Back-action is not Ultra



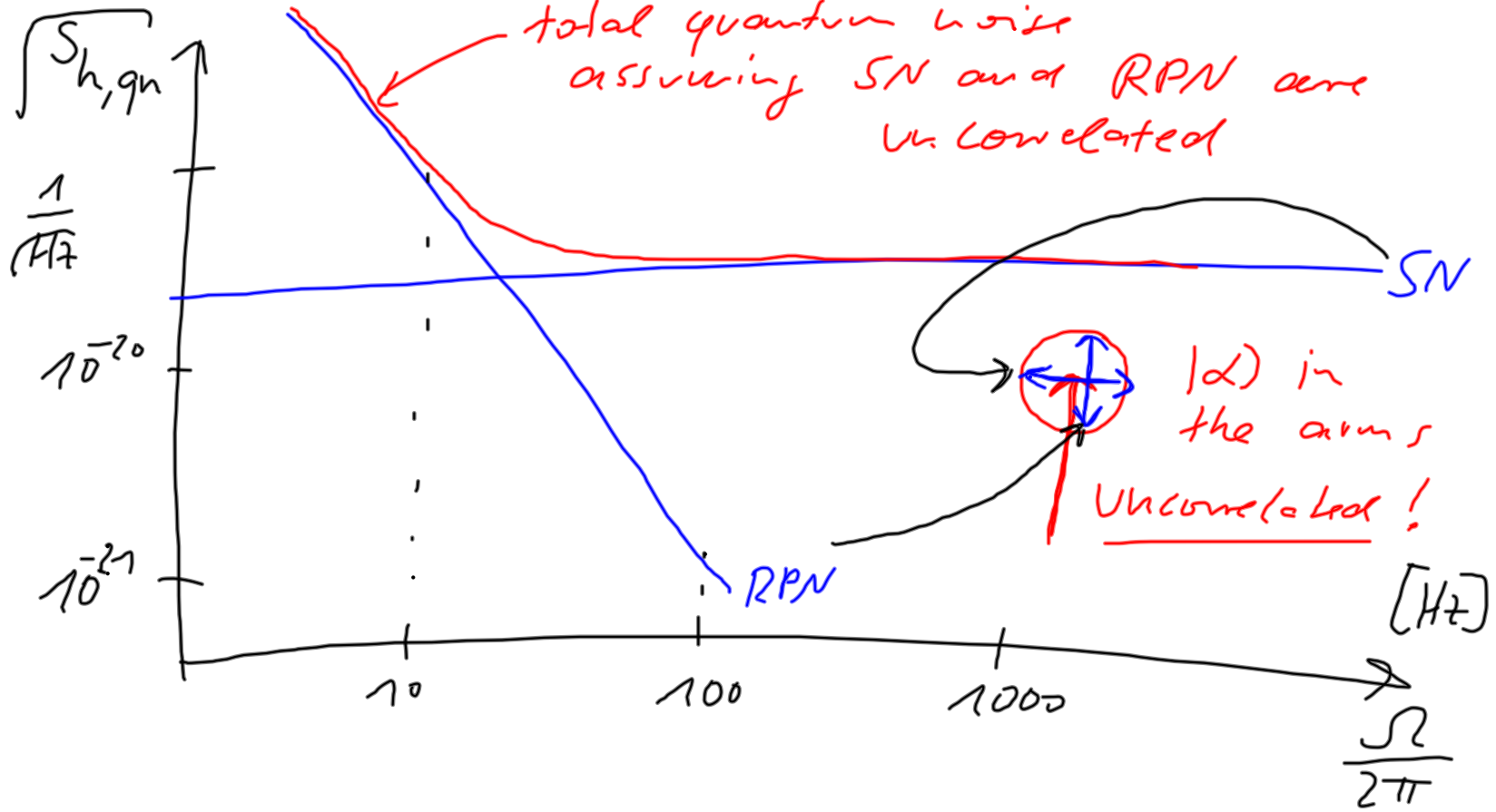
For  $\Omega \gg \Omega_{res}$ .

$$\sqrt{S_{x, rpm}} \approx \frac{1}{\omega \Omega^2} \sqrt{\frac{8t\omega P_0}{c^2}}$$

Recalibration :  $\Delta h = \Delta x \frac{1}{L}$

$$\Rightarrow \sqrt{S_{h, rpm}} = \frac{1}{L \cdot \omega \cdot \Omega^2} \sqrt{\frac{8t\omega P_0}{c^2}} \quad \text{in } \left[ \frac{1}{\sqrt{\text{Hz}}} \right]$$

# GUD Quantum Noise



## The SQL for a force measurement

Assuming uncorrelated SN and RPN

→ add variances of the two  
and find the minimum - when varying  $P_0$

$$\left( \frac{hc^2}{2\omega P_0} + \frac{8\hbar\omega P_0}{c^2 m^2 \Omega^2} \right) \Big|_{\min / P_0}$$

derivative

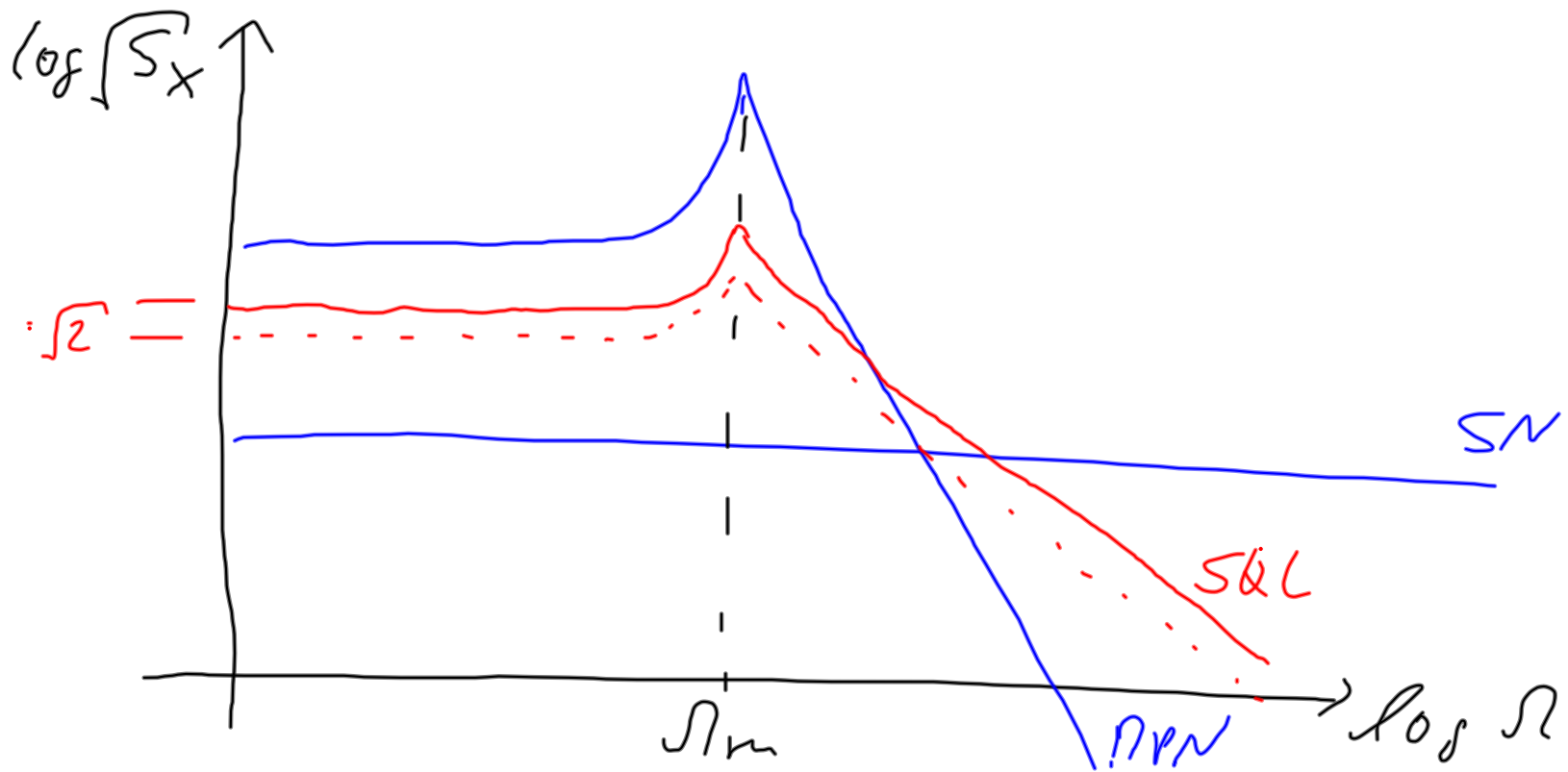
$$\rightarrow P_{0, \text{opt}} = \frac{c^2 m \Omega^2}{4\omega}$$

depends on  
frequency

$$c = \frac{2 \cdot \omega}{2\pi} \quad (\text{over wave length})$$

$$\rightarrow \boxed{S_{X, S\&L} = \frac{4t}{\ln \Omega^2}} \quad \hat{=} \quad S_{X, SN} = S_{X, \Omega PN} = \frac{1}{2} S_{X, S\&L}$$

Spectrum of S&L



The SBL can be suppressed [Jaekel & R. 1915]  
using squeezed light having a  
frequency dependent squeezing angle.

Backaction can be evaded completely  
by using output filter cavities  
"variational output".

The latter is linked to pa den uster  
sweziy.

see Powerpoint slide

—end—